THERMODYNAMICS OF ALCOHOL SOLUTIONS. PHASE EQUILIBRIA AND EXCESS MOLAR ENTHALPIES OF MIXTURES CONTAINING TWO ALCOHOLS

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ABSTRACT

The association constant and the enthalpy of hydrogen-bond formation for pure aliphatic alcohols have been determined from the vapor pressures of the pure alcohols as a function of temperature using Brandani's method, based on the homomorph concept which uses ethers instead of saturated hydrocarbons as the homomorphs of alcohols. The results of phase equilibrium and excess enthalpy data reduction obtained with the new association parameters are comparable with those derived from Brandani's values.

INTRODUCTION

A new method, based on an association model, was presented by Brandani [l] for determining the enthalpy of hydrogen-bond formation and the equilibrium association constant for pure liquid compounds. The method needed the experimental vapor pressures of the pure liquids as a function of temperature and adopted the saturated hydrocarbons as the homomorphs of pure hydrogen-bonded liquids. Brandani's association parameters have been used in the UNIQUAC associated-solution theory to reduce vapor-liquid and liquid-liquid equilibrium and excess molar enthalpy data for associated solutions [2-61. An alternative method [7], based on the Flory-Huggins relation, used the isomeric ether as the alcohol's homomorph and assumed the volume change of mixing. Association parameters derived from the latter method should not be used in the UNIQUAC associated-solution theory, because the theory assumes no volume change.

This paper presents the association constant and the enthalpy of hydrogen-bond formation for pure aliphatic alcohols derived from Brandani's procedure, which adopts the ether as the alcohol's homomorph. The new association parameters will be used with the revised UNIQUAC associatedsolution theory $[4-6]$ to analyze vapor-liquid and liquid-liquid equilibrium and excess molar enthalpy data for mixtures including two alcohols.

ESTIMATION OF ASSOCIATION PARAMETERS FOR ALIPHATIC ALCOHOLS

I assume that the pure alcohol forms linear polymers by consecutive reactions expressed as $A_1 + A_1 = A_{i+1}$ and the equilibrium association constant K_A for these reactions is defined by

$$
K_{A} = \frac{\Phi_{A_{i+1}}^{0}}{\Phi_{A_{i}}^{0} \Phi_{A_{i}}^{0}} \frac{i}{i+1} \text{ for all } i \geq 1
$$
 (1)

 K_A is independent of the degree of association. The enthalpy of hydrogenbond formation h_A fixes the temperature dependence of the association constant by the van't Hoff relation and is assumed to be temperature-independent.

According to Brandani's method [l], the vapor pressure of the pure alcohol is calculated by

$$
P_{A}^{s} = \gamma_{A_{1}}^{0} x_{A_{1}}^{0} P_{h}^{s} \exp[(P_{A}^{s} - P_{h}^{s})(v_{A}^{L} - B_{A}^{F})/RT] \tag{2}
$$

where $\gamma_{\rm A}^{0}$ is the activity coefficient of the pure alcohol monomer adjusted to zero pressure, $x_{A_1}^0$ is the mole fraction of the monomer in the pure alcohol P_h^s is the vapor pressure of the hypothetical fluid, v_A^L is the molar pure alcohol volume estimated from the modified Rackett equation [8] and B_{λ}^{F} is the free contribution to the second virial coefficient calculated from the correlation of Hayden and O'Connell [9] with the related parameters given by Prausnitz et al. [lo].

 $\gamma_{A_1}^0$ is expressed as

$$
\ln \gamma_{A_1}^0 = \ln \frac{\Phi_{A_1}^0}{x_{A_1}^0} + 1 - \frac{\Phi_{A_1}^0}{x_{A_1}^0}
$$
 (3)

where the segment fraction of the pure alcohol monomer ($\Phi_{A_1}^0$) and $x_{A_1}^0$ are given as follows

$$
\Phi_{A_1}^0 = \frac{1 + 2K_A - (1 + 4K_A)^{1/2}}{2K_A^2} \tag{4}
$$

$$
x_{A_1}^0 = 1 - K_A \Phi_{A_1}^0
$$
 (5)

The vapor pressures of the ethers are used to estimate P_h^s . Ambrose et al. [11] have measured the vapor pressures of some ethers as the homomorph of the alcohol: dimethyl ether for ethanol; ethyl methyl ether for I-propanol; propyl methyl ether for 1-butanol; isopropyl methyl ether for isobutanol; butyl methyl ether for 1-pentanol; decyl methyl ether for 1-undecanol. Ambrose et al. have given a general equation to represent their measured vapor pressure data for a number of ethers in terms of the effective carbon number in the ether n^*

$$
\log_{10} P (\text{kPa}) = 7.1972 + 0.1752n^* - (916.74 + 184.766n^*)T^{-1}
$$

-(9.6590 + 1.17110n^*)10⁻⁴T
-(1.34082 - 0.152687n^*)10⁻⁶T² (6)

Substance	Temperature range $(^{\circ}C)$	Absolute arithmetic mean deviation (Torr)	
Methanol	$1.7 - 60$	0.89	
Ethanol	$19.6 - 75$	0.58	
1-Propanol	$48.1 - 90$	0.09	
2-Propanol	$52.3 - 80$	0.06	
1-Butanol	78.6–115	0.18	
2-Butanol	$67.7 - 95$	0.25	
iso-Butanol	$69.9 - 105$	0.24	
tert-Butanol	$56.8 - 80$	0.23	
1-Pentanol	74.8–135	0.66	
1-Hexanol	$52.2 - 155$	1.01	
1-Heptanol	$63.6 - 175$	1.04	
1-Octanol	78.9–190	1.04	
1-Nonanol	$140.0 - 195$	0.69	
1-Decanol	$170.0 - 230$	0.55	

TABLE 1 Results of the fit for vapor pressure data of alcohols

and stated that if used with the values of n^* in table 4 of ref. 11, eqn. (6) reproduces the values obtained in the correlation for all the other compounds, including those whose properties were estimated, nearly as well.

The following values of n^* for ethers corresponding to other alcohols were used: 1 for methanol; 2.60 for 2-propanol, 3.605 for 2-butanol and 3.366 for tert-butanol are halves of those for diisopropyl ether, di-s-butyl ether and di-t-butyl ether, respectively. Values of n^* for 1-hexanol to l-decanol were obtained by linear interpolation between 5.260 of butyl methyl ether for 1-pentanol and 11.3 of decyl methyl ether for 1-undecanol [ll]: 6.267 for 1-hexanol; 7.273 for 1-heptanol; 8.280 for 1-octanol; 9.287 for 1-nonanol; 10.293 for I-decanol.

 K_A at 50°C and h_A for 14 aliphatic alcohols were calculated by minimizing the objective function defined by

$$
F = \sum_{i=1}^{N} [P_{A_i}^s(\exp) - P_{A_i}^s(\text{calc})]^2
$$
 (7)

where eqn. (2) was used to calculate $P_{\rm A}^{\rm s}$ (calc) and $P_{\rm A}^{\rm s}$ (exp) was obtained from the Antoine equation [12] at each of 21 equally spaced points in the temperature range listed in Table 1, where the absolute average deviation between calculated and experimental vapor pressures are reported for 14 alcohols. Table 2 compares the association parameters obtained in this work with those estimated by Brandani [1]. The values of the enthalpy of hydrogen-bond formation obtained here are usually lower than Brandani's values

Substance	K_A (50°C)		$-h_A$ (kJ mol ⁻¹)		
	I ^a	II _p		П	
Methanol	173.9	125.1	25.4	23.6	
Ethanol	110.4	103.2	25.9	23.6	
1-Propanol	87.0	89.9	27.0	23.6	
2-Propanol	49.1	71.9	26.4	23.9	
1-Butanol	69.5	83.0	27.1	23.3	
2-Butanol	31.1	57.7	25.9	22.6	
iso-Butanol	50.6	67.0	26.0	23.0	
tert-Butanol	23.1	30.5	25.9	21.5	
1-Pentanol	47.7	67.8	24.5	22.8	
1-Hexanol	40.8	59.6	24.9	22.4	
1-Heptanol	40.7	49.8	25.9	22.1	
1-Octanol	39.5	41.2	24.9	21.9	
1-Nonanol		32.1		21.9	
1-Decanol	36.8	25.3	20.3	21.8	

TABLE 2

Association parameters for alcohols

^a I. Brandani [1].

b II. This work.

and are comparable to the enthalpies of dilution of ethanol in saturated hydrocarbons at 25° C [13].

ASSOCIATION THEORY

The UNIQUAC associated-solution theory [4-61 gives the activity coefficient of any component in a mixture including two alcohols and the excess molar enthalpy of the mixture as follows.

Binary mixtures containing two alcohols

A and B stand for alcohols. The alcohol B associates linearly as does alcohol A. These two alcohols solvate each other to yield linear copolymers whose general formulas are $(A, B_i)_k$, $A_i(B, A_k)_i$, $(B_i A_i)_k$ and $B_i(A, B_k)_i$, where the indices i , j , k and l can be any integers from 1 to infinity. I assume that the solvated species $(A, B)_{k}$ is not identical with $(B, A_{k})_{k}$, because the last molecule of the former B has a free hydroxyl group and the last one of the latter A has it.

The activity coefficients of both components are expressed by

$$
\ln \gamma_{A} = \ln \left(\frac{\Phi_{A_1}}{\Phi_{A_1}^0 x_A} \right) + r_A \left(\frac{1}{V_A^0} - \frac{1}{V} \right) - \left(\frac{Z}{2} \right) q_A \left(\ln \frac{\Phi_A}{\theta_A} + 1 - \frac{\Phi_A}{\theta_A} \right)
$$

$$
+ q_A \left[-\ln(\theta_A + \theta_B \tau_{BA}) + \theta_B \left(\frac{\tau_{BA}}{\theta_A + \theta_B \tau_{BA}} - \frac{\tau_{AB}}{\theta_B + \theta_A \tau_{AB}} \right) \right]
$$
(8)

$$
\ln \gamma_{\rm B} = \ln \left(\frac{\Phi_{\rm B_{1}}}{\Phi_{\rm B_{1}}^{0} \chi_{\rm B}} \right) + r_{\rm B} \left(\frac{1}{V_{\rm B}^{0}} - \frac{1}{V} \right) - \left(\frac{Z}{2} \right) q_{\rm B} \left(\ln \frac{\Phi_{\rm B}}{\theta_{\rm B}} + 1 - \frac{\Phi_{\rm B}}{\theta_{\rm B}} \right)
$$

$$
+ q_{\rm B} \left[-\ln(\theta_{\rm B} + \theta_{\rm A} \tau_{\rm AB}) + \theta_{\rm A} \left(\frac{\tau_{\rm AB}}{\theta_{\rm B} + \theta_{\rm A} \tau_{\rm AB}} - \frac{\tau_{\rm BA}}{\theta_{\rm A} + \theta_{\rm B} \tau_{\rm BA}} \right) \right]
$$
(9)

where the nominal segment fraction Φ , the nominal surface fraction θ , and the coefficient τ are given by

$$
\Phi_{\mathbf{A}} = r_{\mathbf{A}} x_{\mathbf{A}} / (r_{\mathbf{A}} x_{\mathbf{A}} + r_{\mathbf{B}} x_{\mathbf{B}})
$$
\n
$$
\Phi_{\mathbf{B}} = r_{\mathbf{B}} x_{\mathbf{B}} / (r_{\mathbf{A}} x_{\mathbf{A}} + r_{\mathbf{B}} x_{\mathbf{B}})
$$
\n
$$
\theta_{\mathbf{A}} = q_{\mathbf{A}} x_{\mathbf{A}} / (q_{\mathbf{A}} x_{\mathbf{A}} + q_{\mathbf{B}} x_{\mathbf{B}})
$$
\n(10)

$$
\theta_{\mathbf{B}} = q_{\mathbf{B}} \mathbf{x}_{\mathbf{B}} / (q_{\mathbf{A}} \mathbf{x}_{\mathbf{A}} + q_{\mathbf{B}} \mathbf{x}_{\mathbf{B}})
$$
(11)

$$
\tau_{BA} = \exp(-a_{BA}/I)
$$

\n
$$
\tau_{AB} = \exp(-a_{AB}/T)
$$
\n(12)

^rand q are the pure-component structural parameters estimated by the method of Vera et al. [14] and the energy parameters, a_{BA} and a_{AB} , are obtained in fitting the theory to experimental phase equilibrium data.

The monomer segments fractions, Φ_{A_1} and Φ_{B_1} , are obtained from the following mass balance equations

$$
\Phi_{\mathbf{A}} = \overline{S}_{\mathbf{A}} + \frac{r_{\mathbf{A}} K_{\mathbf{A}\mathbf{B}} \overline{S}_{\mathbf{A}} S_{\mathbf{B}}}{\left(1 - r_{\mathbf{A}} r_{\mathbf{B}} K_{\mathbf{A}\mathbf{B}}^2 S_{\mathbf{A}} S_{\mathbf{B}}\right)^2} \times \left[2 + r_{\mathbf{B}} K_{\mathbf{A}\mathbf{B}} S_{\mathbf{A}} \left(2 - r_{\mathbf{A}} r_{\mathbf{B}} K_{\mathbf{A}\mathbf{B}}^2 S_{\mathbf{A}} S_{\mathbf{B}}\right) + r_{\mathbf{A}} K_{\mathbf{A}\mathbf{B}} S_{\mathbf{B}}\right]
$$
\n(13)

$$
\Phi_{\rm B} = \overline{S}_{\rm B} + \frac{r_{\rm B} K_{\rm AB} S_{\rm A} \overline{S}_{\rm B}}{\left(1 - r_{\rm A} r_{\rm B} K_{\rm AB}^2 S_{\rm A} S_{\rm B}\right)^2}
$$
\n
$$
\sqrt{2 + r_{\rm B} K_{\rm AB}^2 S_{\rm A} S_{\rm B}} = (14)
$$

$$
\times \left[2 + r_A K_{AB} S_B (2 - r_A r_B K_{AB}^2 S_A S_B) + r_B K_{AB} S_A\right]
$$
(14)

where the sums S_A , S_B , S_A and S_B are defined by

$$
\overline{S}_{A} = \Phi_{A_{1}} / (1 - K_{A} \Phi_{A_{1}})^{2}
$$
 (15)

$$
\overline{S}_{\mathbf{B}} = \Phi_{\mathbf{B}_1} / (1 - K_{\mathbf{B}} \Phi_{\mathbf{B}_1})^2
$$
 (16)

$$
S_{A} = \Phi_{A_1} / (1 - K_A \Phi_{A_1})
$$
\n⁽¹⁷⁾

$$
S_{\mathbf{B}} = \Phi_{\mathbf{B}_1} / (1 - K_{\mathbf{B}} \Phi_{\mathbf{B}_1})
$$
\n(18)

The true molar volume of the binary mixture V is expressed as

$$
\frac{1}{V} = \frac{S_A}{r_A} + \left(\frac{2}{r_A r_B K_{AB}} + \frac{S_A}{r_A} + \frac{S_B}{r_B}\right) \frac{r_A r_B K_{AB}^2 S_A S_B}{\left(1 - r_A r_B K_{AB}^2 S_A S_B\right)} + \frac{S_B}{r_B}
$$
(19)

At pure alcohol states, $\Phi_{A_1}^0$ is given by eqn. (4). $\Phi_{B_1}^0$, V_A^0 and V_B^0 are obtained bY

$$
\Phi_{\mathbf{B}_1}^0 = \frac{1 + 2K_{\mathbf{B}} - (1 + 4K_{\mathbf{B}})^{1/2}}{2K_{\mathbf{B}}^2}
$$
(20)

$$
1/V_{\rm A}^0 = \left(1 - K_{\rm A} \Phi_{\rm A_1}^0\right)/r_{\rm A} \tag{21}
$$

$$
1/V_{\rm B}^0 = (1 - K_{\rm B} \Phi_{\rm B_1}^0)/r_{\rm B}
$$
 (22)

The excess molar enthalpy of the mixture is given by

$$
H^{E} = h_{A}x_{A}\left(\frac{\overline{U}_{A}\Phi_{A_{1}}}{\Phi_{A}} - \overline{U}_{A}^{0}\Phi_{A_{1}}^{0}\right) + h_{B}x_{B}\left(\frac{\overline{U}_{B}\Phi_{B_{1}}}{\Phi_{B}} - \overline{U}_{B}^{0}\Phi_{B_{1}}^{0}\right)
$$

+
$$
\left\{h_{A}\left[\frac{\overline{U}_{A}}{K_{AB}U_{A}}\left(\frac{x_{B}}{\Phi_{B}r_{A}} + \frac{x_{A}}{\Phi_{A}r_{B}}\right) + \frac{\overline{U}_{A}x_{A}\Phi_{A_{1}}}{\Phi_{A}}\left(2 - r_{A}r_{B}K_{AB}^{2}\Phi_{A}\Phi_{B_{1}}U_{A}U_{B}\right)\right]
$$

+
$$
\frac{\overline{U}_{A}U_{B}x_{B}\Phi_{B_{1}}}{U_{A}\Phi_{B}}\right] + h_{B}\left[\frac{\overline{U}_{B}}{K_{AB}U_{B}}\left(\frac{x_{B}}{\Phi_{B}r_{A}} + \frac{x_{A}}{\Phi_{A}r_{B}}\right)\right]
$$

+
$$
\frac{\overline{U}_{B}x_{B}\Phi_{B_{1}}}{\Phi_{B}}\left(2 - r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}\right) + \frac{U_{A}\overline{U}_{B}x_{A}\Phi_{A_{1}}}{U_{B}\Phi_{A}}\right]
$$

+
$$
h_{AB}\left[\left(\frac{x_{B}}{\Phi_{B}r_{A}} + \frac{x_{A}}{\Phi_{A}r_{B}}\right)\frac{\left(1 + r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}\right)}{K_{AB}}\right]
$$

+
$$
2\left(\frac{U_{A}x_{A}\Phi_{A_{1}}}{\Phi_{A}} + \frac{U_{B}x_{B}\Phi_{B_{1}}}{\Phi_{B}}\right)\right]\right\}\frac{r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}}{(1 - r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B})^{2}}
$$

-
$$
R\left[\frac{q_{A}x_{A}\theta_{B}}{
$$

where \overline{U}_A , \overline{U}_B , U_A and U_B , respectively, are defined by

$$
\overline{U}_{A} = K_{A} \Phi_{A_{1}} / (1 - K_{A} \Phi_{A_{1}})^{2}
$$
\n(24)

$$
\overline{U}_{\rm B} = K_{\rm B} \Phi_{\rm B_1} / (1 - K_{\rm B} \Phi_{\rm B_1})^2
$$
 (25)

$$
U_{\mathsf{A}} = 1/(1 - K_{\mathsf{A}} \Phi_{\mathsf{A}_1}) \tag{26}
$$

$$
U_{\rm B} = 1/(1 - K_{\rm B} \Phi_{\rm B_1})\tag{27}
$$

It is assumed that a_{BA} and a_{AB} are linearly dependent on temperature. $a_{BA} = C_A + D_A (T - 273.15)$ $a_{AB} = C_B + D_B(T - 273.15)$ (28)

Ternary mixtures containing two alcohols and one active nonassociating component

In the mixture exist associated species and solvated ones whose general forms are A_t, B_t, $(A_tB_j)_k$, $A_t(B_jA_k)_t$, $(B_tA_j)_k$, $B_t(A_jB_k)_t$, $A_tC_t(A_iB_j)_kC_t$ $A_{\lambda}(B, A_{\lambda})$, C, B, C, $(B_{\lambda}A_{\lambda})$ _kC and $B_{\lambda}(A, B_{\lambda})$, C, where the indices *i*, *j*, *k* and *l* can go from 1 to infinity.

The activity coefficient of component I is

$$
\ln \gamma_I = \ln \left(\frac{\Phi_{I_1}}{\Phi_{I_1}^0 x_I} \right) + \frac{r_I}{V_I^0} - \frac{r_I}{V} - \left(\frac{Z}{2} \right) q_I \left(\ln \frac{\Phi_I}{\Phi_I} + 1 - \frac{\Phi_I}{\Phi_I} \right)
$$

+
$$
q_I \left[1 - \ln \left(\sum_J \theta_I \tau_{JI} \right) - \sum_J \frac{\theta_J \tau_{IJ}}{\sum_K \theta_K \tau_{KJ}} \right]
$$
(29)

where Φ_i , θ_i , V^0 , V and τ_{ij} are expressed as follows:

$$
\Phi_I = r_I x_I / \sum_J r_J x_J \tag{30}
$$

$$
\theta_I = q_I x_I / \sum_J q_J x_J \tag{31}
$$

$$
1/V_I^0 = (1 - K_I \Phi_{I_1}^0)/r_I
$$
\n(32)

$$
\frac{1}{V} = \frac{S_A}{r_A} + \frac{S_B}{r_B} + \left(\frac{2}{r_A r_B K_{AB}} + \frac{S_A}{r_A} + \frac{S_B}{r_B}\right) \frac{r_A r_B K_{AB} S_A S_B}{(1 - r_A r_B K_{AB}^2 S_A S_B)}
$$
\n
$$
+ \frac{\Phi_{C_1}}{r_C} \left\{1 + r_C K_{AC} S_A + r_C K_{BC} S_B + \left(\frac{1}{r_A K_{AB}} + S_B\right) K_{BC}\right\} \frac{r_A r_B r_C K_{AB}^2 S_A S_B}{(1 - r_A r_B K_{AB}^2 S_A S_B)}\right\}
$$
\n(33)

$$
\tau_{JI} = \exp(-a_{JI}/T)
$$
\n
$$
\Phi_{C}^{0} = 1 \text{ for component C, } \Phi_{A}^{0} \text{ and } \Phi_{B}^{0} \text{ are given by eqns. (4) and (20).}
$$
\n(34)

The monomer segment fractions, Φ_{A_1} , Φ_{B_1} and Φ_{C_1} , are numerical solved from the following mass balance equations:

$$
\Phi_{A} = (1 + r_{A} K_{AC} \Phi_{C_{1}}) \bar{S}_{A} + \frac{r_{A} K_{AB} S_{A} S_{B}}{(1 - r_{A} r_{B} K_{AB}^{2} S_{A} S_{B})^{2}} \times \left\{ 2 + r_{B} K_{AB} S_{A} (2 - r_{A} r_{B} K_{AB}^{2} S_{A} S_{B}) + r_{A} K_{AB} S_{B} \n+ \Phi_{C_{1}} \left[(r_{A} K_{AC} + r_{B} K_{BC}) + r_{A} r_{B} K_{AB} K_{AC} S_{A} \n\times (2 - r_{A} r_{B} K_{AB}^{2} S_{A} S_{B}) + r_{A} r_{B} K_{AB} K_{BC} S_{B} \right] \right\}
$$
\n(35)

$$
\Phi_{B} = (1 + r_{B} K_{BC} \Phi_{C_{1}}) \overline{S}_{B} + \frac{r_{B} K_{AB} S_{A} \overline{S}_{B}}{(1 - r_{A} r_{B} K_{AB}^{2} S_{A} S_{B})^{2}} \times (2 + r_{A} K_{AB} S_{B} (2 - r_{A} r_{B} K_{AB}^{2} S_{A} S_{B}) + r_{B} K_{AB} S_{A} + \Phi_{C_{1}} [(r_{A} K_{AC} + r_{B} K_{BC}) + r_{A} r_{B} K_{AB} K_{BC} S_{B} \times (2 - r_{A} r_{B} K_{AB} S_{A} S_{B}) + r_{A} r_{B} K_{AB} K_{AC} S_{A}]
$$
\n(36)

$$
\Phi_{C} = \Phi_{C_{1}} \left\{ 1 + r_{C} K_{AC} S_{A} + r_{C} K_{BC} S_{B} + \frac{r_{A} r_{B} r_{C} K_{AB}^{2} S_{A} S_{B}}{\left(1 - r_{A} r_{B} K_{AB}^{2} S_{A} S_{B} \right)} \left[\frac{K_{AC}}{r_{B} K_{AB}} + \frac{K_{BC}}{r_{A} K_{AB}} + K_{AC} S_{A} + K_{BC} S_{B} \right] \right\}
$$
(37)

The excess molar enthalpy of the mixture is expressed as

$$
H^{E} = h_{A}x_{A}\left(\frac{\overline{U}_{A}\Phi_{A_{1}}}{\Phi_{A}} - \overline{U}_{A}^{0}\Phi_{A_{1}}^{0}\right) + h_{B}x_{B}\left(\frac{\overline{U}_{B}\Phi_{B_{1}}}{\Phi_{B}} - \overline{U}_{B}^{0}\Phi_{B_{1}}^{0}\right) + (h_{A}\overline{U}_{A} + h_{AC}U_{A})\frac{r_{A}K_{AC}\Phi_{C_{1}X_{A}\Phi_{A_{1}}}}{\Phi_{A}} + (h_{B}\overline{U}_{B} + h_{BC}U_{B})\frac{r_{B}K_{BC}\Phi_{C_{1}X_{B}\Phi_{B_{1}}}}{\Phi_{B}} + \left(h_{A}\left(\frac{\overline{U}_{A}}{K_{AB}U_{A}}\left(\frac{x_{B}}{r_{A}\Phi_{B}} + \frac{x_{A}}{r_{B}\Phi_{A}}\right) + \frac{\overline{U}_{A}x_{A}\Phi_{A_{1}}}{\Phi_{A}}\left(2 - r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}\right)\right) + \frac{\overline{U}_{A}U_{B}x_{B}\Phi_{B_{1}}}{U_{A}\Phi_{B}} + \Phi_{C_{1}}\left[\left(\frac{r_{B}K_{BC}x_{B}}{r_{A}K_{AB}\Phi_{B}} + \frac{r_{A}K_{AC}x_{A}}{r_{B}K_{AB}\Phi_{A}}\right)\frac{\overline{U}_{A}}{U_{A}} + \frac{r_{A}K_{AC}\overline{U}_{A}x_{A}\Phi_{A_{1}}}{\Phi_{A}}\left(2 - r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}\right) + \frac{r_{B}K_{BC}\overline{U}_{A}U_{B}x_{B}\Phi_{B_{1}}}{U_{A}\Phi_{B}}\right]\right) + h_{B}\left(\frac{\overline{U}_{B}}{K_{AB}U_{B}}\left(\frac{x_{B}}{r_{A}\Phi_{B}} + \frac{x_{A}}{r_{B}\Phi_{A}}\right) + \frac{\overline{U}_{B}x_{B}\Phi_{B_{1}}}{\Phi_{B}}\left(2 - r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}\right) \right) + \frac{U_{A
$$

$$
+\Phi_{C_{1}}\left[\left(\frac{r_{B}K_{BC}x_{B}}{r_{A}\Phi_{B}} + \frac{r_{A}K_{AC}x_{A}}{r_{B}\Phi_{A}}\right) \times \frac{\left(1 + r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}\right)}{K_{AB}} + 2\left(\frac{r_{A}K_{AC}U_{A}x_{A}\Phi_{A_{1}}}{\Phi_{A}} + \frac{r_{B}K_{BC}U_{B}x_{B}\Phi_{B_{1}}}{\Phi_{B}}\right)\right]\right) + \left[h_{AC}r_{A}K_{AC}\Phi_{C_{1}}\left(\frac{U_{A}x_{A}\Phi_{A_{1}}}{\Phi_{A}} + \frac{x_{A}}{r_{A}K_{AB}\Phi_{A}}\right) + h_{BC}r_{B}K_{BC}\Phi_{C_{1}}\times \left(\frac{U_{B}x_{B}\Phi_{B_{1}}}{\Phi_{B}} + \frac{x_{B}}{r_{A}K_{AB}\Phi_{B}}\right)\right]\left(1 - r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}\right)\right) \times \frac{r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}}{\left(1 - r_{A}r_{B}K_{AB}^{2}\Phi_{A_{1}}\Phi_{B_{1}}U_{A}U_{B}\right)^{2}} - R\sum_{i}q_{i}x_{i} \frac{\sum_{j} \theta_{j} \frac{\partial \tau_{ij}}{\partial(1/T)}}{\sum_{j} \theta_{j} \tau_{ji}} \tag{38}
$$

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(39)

where a_{II} is expressed by $a_{11} = C_1 + D_1(T-273.15)$

CALCULATED RESULTS

Binary mixtures

The excess molar enthalpies of binary mixtures containing normal aliphatic alcohols are positive (heat is absorbed) and those of some binary mixtures of normal and branched alcohols are negative (heat is evolved). The theory is able to reproduce these data as shown below. Table 3 shows the binary calculated results for alcohol-alcohol mixtures obtained by use of several sets of K_{AB} and h_{AB} without addition of the energy parameters and some of the results are compared with experimental values in Figs. 1 and 2. These figures indicate that the experimental data of methanol-2-propanol and ethanol-1-pentanol mixtures are reproduced with only selected sets of K_{AB} and h_{AB} and this is not the case for methanol-1-butanol, where additional energy parameters should be included for the accurate correlation of the experimental data.

Vapor-liquid equilibrium data reduction was carried out by use of the thermodynamic equation

$$
\phi_I y_I P = \gamma_I x_I \phi_I^s P_I^s \exp\left[v_I^L (P - P_I^s) / RT\right]
$$
\n(40)

where y , P , P^s and R are the vapor-phase mole fraction, total pressure,

TABLE 3

Results of fitting the UNIQUAC associated-solution theory to excess molar enthalpies for binary alcohol-alcohol mixtures at 25° C with the chemical contribution term alone

pure-component vapor pressure and gas constant. The fugacity coefficients, ϕ_t and ϕ_t^s , were calculated from the volume-explicit equation of state truncated after the second term. The second virial coefficients were estimated from the method of Hayden and O'Connell [9]. Pure-component vapor pressures were calculated from the Antoine equation whose parameters were taken from the literature [20,21].

The computer program used, based on the maximum likelihood principle, was similar to that described by Prausnitz et al. [10]. The standard deviations in the measured variables were taken as $\sigma_p = 1.0$ Torr, $\sigma_T = 0.05$ K, $\sigma_{\nu} = 0.001$ and $\sigma_{\nu} = 0.003$.

Parameter estimation from mutual solubilities was performed by solving eqn. (41) for any component I .

$$
\left(x_{I}\gamma_{I}\right)^{I}=\left(x_{I}\gamma_{I}\right)^{II}
$$
\n(41)

where the indices, I and II, represent equilibrium liquid phases.

Table 4 lists the values of the solvation equilibrium constants and enthalpies of complex formation used in binary and ternary phase equilibrium and excess enthalpy data calculations.

Table 5 presents the calculated result of binary phase equilibrium data reduction. In the correlation of vapor-liquid equilibria for alcohol-alcohol mixtures, various values of K_{AB} were used with or without the energy parameters. For a fixed value of K_{AR} the results obtained with the energy parameters are better than those without the energy parameters. Many sets of values of K_{AR} and the energy parameters can reduce the experimental data of these mixtures with good accuracy. Table 6 gives the calculated

Fig. 1. Excess molar enthalpies for (a) methanol $(1)-2$ -propanol (2) and (b) ethanol $(1)-1$ pentanol (2) at 25°C. Experimental (\bullet): (a) ref. 16; (b) ref. 17. Calculated: (a) $\left(-\cdots\right)$ $K_{AB} = 70(50^{\circ} \text{C})$ and $h_{AB} = -23.75 \text{ kJ} \text{ mol}^{-1}$; (**kJ mole'. (______) KA8-** -) KA, = 76(5O"C) and *h,, = -* 23.75 $A_{\rm B} = 80(50^{\circ}\text{C})$ and $h_{\rm AB} = -23.75$ kJ mol⁻¹; (b) (------) $K_{\rm AB} =$ 33(50°C) and $h_{AB} = -23.2$ kJ mol⁻¹; (-----) $K_{AB} = 43(50^{\circ}C)$ and $h_{AB} = -23$ kJ mol⁻¹ $(- - - -)$ $K_{AB} = 36(50^{\circ}\text{C})$ and $h_{AB} = -23.2 \text{ kJ}$ mol

results of excess molar enthalpy. The accuracy of the correlation studied in this paper is the same as that obtained in previous papers $[4-6]$, indicating that the results of data reduction are not sensitive to used association parameters.

Fig. 2. Excess molar enthalpies for methanol (1)-1-butanol (2) at 25 $^{\circ}$ C. Experimental (0), ref. 31. Calculated: $(- - - -)$ $K_{AB} = 46(50^{\circ}C)$ and $h_{AB} = -23.45$ kJ mol⁻¹; $(K_{AB}$ = 56(50°C) and h_{AB} = -23.3 kJ mol⁻¹; (------) K_{AB} = 50(50°C) and h_{AB} = -23.4 kJ mol^{-1} .

TABLE 4

Solvation constants and enthalpies of complex formation

Mixture $(A-B)$	K_{AB} (50°C)	$-h_{AB}$ (kJ mol ⁻¹)	
Methanol-benzene		8.3	
Methanol-chlorobenzene	$3(55^{\circ}C)$		
Methanol-ethanol	110	23.6	
Methanol-2-propanol	80	23.75	
Ethanol-benzene	2	8.3	
2-Propanol-chlorobenzene	$2(55^{\circ}C)$		
2-Propanol-benzene	1.8	8.3	

 $^{\rm a}$ Parameters used for ternary phase equilibrium calculations. $^{\rm b}$ MS = mutual solubilities. ^a Parameters used for ternary phase equilibrium calculations.

 MS = mutual solubilities.

Ternary mixtures

Table 7 shows the ternary predicted results of excess molar enthalpies for 18 mixtures containing two alcohols and one saturated hydrocarbon and one

TABLE 7

Predicted results of excess molar enthalpies for ternary mixtures containing two alcohols and one hydrocarbon at 25°C

TABLE 8

Ternary predicted results of vapor-liquid equilibria obtained from binary parameters

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Fig. 3. Solubility envelopes for mixtures containing two alcohols and one saturated hydrocarbon. Experimental: tie-line (\bullet) ; solubility (O) . Calculated $(____\)$. (A) Methanol-ethanol-cyclohexane at 25° C [39]; (B) methanol-ethanol-methylcyclohexane at 25° C [30]; (C) methanol-2-propanol-n-hexane at 5° C [29]; (D) methanol-2-propanol-n-hexane at 25° C [29]; (E) methanol-2-propanol-cyclohexane at 25° C [40].

mixture including ethanol, 1-propanol and benzene. The overall mean deviation for 19 mixtures is 10.47 J mol⁻¹, which corresponds to 10.59 J mol⁻¹ obtained in previous papers [5,6]. The predicted results of vapor-liquid equilibria for two mixtures are given in Table 7. The magnitude of the deviations is similar to that obtained previously [4]. The values of K_{AB} for alcohol-alcohol mixtures influence the calculated results of ternary liquid-liquid equilibria for mixtures containing two alcohols. The parameter set of $K_{AB} = 80$ at 50°C and non-zero values of the energy parameters for

the methanol-ethanol mixture provided too large ternary solubility envelopes for methanol-ethanol-cyclohexane and methanol-ethanol-methylcylohexane mixtures. The parameter set of $K_{AB} = 110$ at 50°C correctly reproduces liquid-liquid equilibria for these mixtures as shown in Fig. 3. Out of the parameters sets for the methanol-2-propanol mixture, the value of $K_{AB} = 80$ at 50°C gives the best calculated results of ternary liquid-liquid equilibria for methanol-2-propanol-n-hexane and methanol-2-propanolcyclohexane mixtures (Fig. 3).

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LIST OF SYMBOLS

Greek letters

Subscripts

Subscript

0 pure-liquid reference state

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